#### Zhaofeng Kang, KIAS, 20/03/2015

Exploring the Dark Sector/KIAS/Seoul/ March 16 (Mon) – March 20 (Fri), 2015

Outline

- Glory of WIMP dark matter
- Predicament of WIMP dark matter
- Quark-phobic dark matter
- Conclusion

## **Glory of WIMP DM**

### • Well-known for the WIMP miracle

DM annihilates involving weak scale particles (~100 GeV) and weak scale strength (~0.2) naturally with rate~1pb.

- Predicted by supersymmetry!!!
- More generically, connected to TeV new physics,

Such as hierarchy problem, neutrino and flavor physics, etc., servicing as a model source.

Promoting a long research chain, along which —

Particle, astronomical, cosmological and nuclear

### **Predicament of WIMP DM**

#### Direct detection attacks

#### Limits from LUX (heavy>mb) & CDMSlite (light<mb)



Anomalies from GoGeNT, etc., ~10 GeV DM? ? Isospin-violating DM to reconcile them with LUX, etc.??? An inspiring coincidence:

$$\frac{\langle \sigma_{an} v \rangle}{\sigma_p} \sim \frac{M_{\rm DM}^2}{\mu_p^2}, \quad \mu_p \simeq 1 \, {\rm GeV}. \label{eq:sigma_angle}$$

## **Predicament of WIMP DM**

### Indirect detection attacks

#### AMS-02 (positron killer) & Fermi-LAT (Gamma killer)



Ice cube (neutrino killer) & Planck CMB (sensitive to *p*-wave annihilating DM!!!)



The light DM (100 GeV) with

S-Wave annihilation is in danger

### **Predicament of WIMP DM**

Implications to DM with SM mediators

#### DM with Z-portal:

DM with Higgs portal

$$\frac{1}{2}$$
The Higgs portal
$$\frac{1}{12} = \frac{1}{12} \cdot \frac{1}$$

for the scalar DM,

λ~λv/2mom 510-2 50

Thus, they can not account for correct relic density for DM. The WIMP miracle requires new mediators. Or, we give up the miracle. Anyway, we can not jump to the conclusion that WIMP DM encounters a real crisis ---

### Actually, there are a lot of ways to break siege

Tuning scenarios: Annihilating with a resonance (near Higgs pole) or coannihilating with a partner (degenerate LSP & stop in SUSY) Point: allow very weakly Axial mediators: A fermionic DM with CP-odd

Higgs or axial-vector coupling, etc.,

Only gives spin-dependent **DM-nucleon scattering** 



coupling between DM

and the mediators

Quark-phobic DM?

DM annihilates into the dark states, which feebly couple to SM

one dead, thousands born!!

Easy ways to quark-phobic DM

The presence of a new sector, to which DM tightly couples to. The new sector does not lead to DM coupling to quarks

Easily defeats LUX & superCDMS



Consistent with leptonic DM for XXX-anomaly?

Old famous anomaly: excess in e<sup>+</sup> but not p<sup>+</sup>, Very new anomaly: Galaxy center GeV- $\gamma$  excess Very old anomaly: muon g-2!

#### More on GeV-anomaly

Well fit by thermal DM DM $\rightarrow$ bb/ $\tau\tau$ , tightly constrained by dSphs. The *bb*-mode is in strong tension with dSphs while  $\tau\tau$  is fine!!



The spectrum from Ret2 clearly rises above the expected background from 2 to 10 GeV, thus disfavoring millisecond pulsar but favoring DM:

A support from Reticulum 2? It is (the only) one of the 9 newly discovered dSphs that sees  $\gamma$  excess. In the DM explanation, the best fit is  $\mu\mu$  and the next is  $\tau\tau$ .



### Several known ways (inspired by PAMELA etc.)

DM from a dark sector that has good DM dynamics but feebly couples to SM for thermalizing, via small kinematic/mass mixing N-Arkani-Hamed, D.-P.-Finkbeiner, T.-R.-Slatyer and N.-Weiner, Phys.Rev. D79 (2009) 015014 A new leptonic gauge group like  $U(1)_L$  or  $U(1)_{Li-Lj}$  under which DM is charged, as leads to leptonic DM. P. J. Fox and E. Poppitz, Phys. Rev. D 79, 083528 (2009). S. Baek and P. Ko, JCAP 0910, 011 (2009). A t-channel leptonic mediator to the leptons, e.g., slepton in the supersymmetric standard models

#### Focused example: via a leptonic Higgs doublet —

H. S. Goh, L. J. Hall and P. Kumar, JHEP 0905, 097 (2009)

M. S. Boucenna and S. Profumo, Phys. Rev. D 84, 055011 (2011)

### Ex: lepton-specific 2HDM (S+2HDML)

2HDML is characterized by a lepophilic Higgs doublet  $\Phi_L = \Phi_1$ ,

 $-\mathcal{L}_Y = Y^u \overline{Q_L} \widetilde{\Phi}_2 u_R + Y^d \overline{Q_L} \Phi_2 d_R + Y^e \overline{l_L} \Phi_1 e_R + c.c.,$ 

**Embedded in**  $U(1)_{PQ}$  **models**? *S* is a complex scalar DM stabilized by a remnant of  $U(1)_{PQ}$ ?

 $\Phi_L$  provides a leptonic portal for the DM S, a real singlet scalar:

$$-\mathcal{L}_{\rm DM} = \frac{1}{2}m_S^2 S^2 + \left(\frac{\eta_1}{2}S^2 \Phi_2^{\dagger} \Phi_L + c.c.\right) + \frac{\eta_3}{2}S^2 |\Phi_L|^2 + \frac{\eta_2}{2}S^2 |\Phi_2|^2 + \frac{\eta_4}{4}S^4$$

*h*<sub>SM</sub>-S-S is suppressed by a large tanβ~100, leaving only  $Φ_L$ -portal to ττ, or to *H h*<sub>SM</sub>

For  $\tan\beta >>1$ , only the contact interactions between DM and the leptonic Higgs are sizable In the supersymmetric version, this term, from the F-term, can be naturally suppressed

Je A/H/HL

### • Easy g-2 with a light A & large tanβ

In general, the CP-even Higgs bosons  $h_{SM}$  and H contribute negative to g-2 while A (>GeV) contribute positive — —



The two-loop Zee-Babu diagram tends to be dominant



#### In summary, we expect a lighter A but heavier

 $H \sim H^+$ , to recover the custodial SU(2). They are testable at LHC

In other types of 2DHM like type-II, it is shown that the theoretical and B-physics constraints **negate a light A for g-2**, see 1409.3199, by A. Broggio etc



### Easy leptophilic dark matter

With the leptonic portal  $\Phi_L$ , DM safely annihilates into  $\tau\tau$  via H or directly into the portal, i.e., a pair of Higgs bosons AA/HH...

$$\langle \sigma v \rangle_{\tau\tau} \approx \frac{\eta_1^2}{64\pi} \frac{1}{m_S^2} \left( \frac{m_\tau}{m_S} \right)^2 \frac{\tan^2 \beta}{(1 - m_H^2/4m_S^2)^2} \left( 1 - \frac{m_\tau^2}{m_S^2} \right)^{3/2} \qquad \qquad \langle \sigma v \rangle_{H_iH_i} \approx \frac{1}{64\pi} \frac{\eta_3^2}{m_S^2} C_i \left( 1 - \frac{m_{H_i}^2}{m_S^2} \right)^{3/2} = \frac{1}{1 - 1} \frac{\eta_2^2}{m_S^2} \left( 1 - \frac{m_H^2}{m_S^2} \right)^{3/2} = \frac{1}{1 - 1} \frac{\eta_2^2}{m_S^2} \left( 1 - \frac{m_H^2}{m_S^2} \right)^{3/2} = \frac{1}{1 - 1} \frac{\eta_2^2}{m_S^2} \left( 1 - \frac{m_H^2}{m_S^2} \right)^{3/2} = \frac{1}{1 - 1} \frac{\eta_2^2}{m_S^2} \left( 1 - \frac{m_H^2}{m_S^2} \right)^{3/2} = \frac{1}{1 - 1} \frac{\eta_2^2}{m_S^2} \left( 1 - \frac{m_H^2}{m_S^2} \right)^{3/2} = \frac{1}{1 - 1} \frac{\eta_2^2}{m_S^2} \left( 1 - \frac{m_H^2}{m_S^2} \right)^{3/2} = \frac{1}{1 - 1} \frac{\eta_2^2}{m_S^2} \left( 1 - \frac{\eta_2^2}{m_S^2} \right)^{3/2} = \frac{1}{1 - 1} \frac{\eta_2^2}{m_S^2} \left( 1 - \frac{\eta_2^2}{m_S^2} \right)^{3/2} = \frac{1}{1 - 1} \frac{\eta_2^2}{m_S^2} \left( 1 - \frac{\eta_2^2}{m_S^2} \right)^{3/2} = \frac{1}{1 - 1} \frac{\eta_2^2}{m_S^2} \left( 1 - \frac{\eta_2^2}{m_S^2} \right)^{3/2} = \frac{1}{1 - 1} \frac{\eta_2^2}{m_S^2} \left( 1 - \frac{\eta_2^2}{m_S^2} \right)^{3/2} = \frac{1}{1 - 1} \frac{\eta_2^2}{m_S^2} \left( 1 - \frac{\eta_2^2}{m_S^2} \right)^{3/2} = \frac{1}{1 - 1} \frac{\eta_2^2}{m_S^2} \left( 1 - \frac{\eta_2^2}{m_S^2} \right)^{3/2} = \frac{1}{1 - 1} \frac{\eta_2^2}{m_S^2} \left( 1 - \frac{\eta_2^2}{m_S^2} \right)^{3/2} = \frac{1}{1 - 1} \frac{\eta_2^2}{m_S^2} \left( 1 - \frac{\eta_2^2}{m_S^2} \right)^{3/2} = \frac{1}{1 - 1} \frac{\eta_2^2}{m_S^2} \left( 1 - \frac{\eta_2^2}{m_S^2} \right)^{3/2} = \frac{1}{1 - 1} \frac{\eta_2^2}{m_S^2} \left( 1 - \frac{\eta_2^2}{m_S^2} \right)^{3/2} = \frac{1}{1 - 1} \frac{\eta_2^2}{m_S^2} \left( 1 - \frac{\eta_2^2}{m_S^2} \right)^{3/2} = \frac{1}{1 - 1} \frac{\eta_2^2}{m_S^2} \left( 1 - \frac{\eta_2^2}{m_S^2} \right)^{3/2} = \frac{1}{1 - 1} \frac{\eta_2^2}{m_S^2} \left( 1 - \frac{\eta_2^2}{m_S^2} \right)^{3/2} = \frac{1}{1 - 1} \frac{\eta_2^2}{m_S^2} \left( 1 - \frac{\eta_2^2}{m_S^2} \right)^{3/2} = \frac{1}{1 - 1} \frac{\eta_2^2}{m_S^2} \left( 1 - \frac{\eta_2^2}{m_S^2} \right)^{3/2} = \frac{1}{1 - 1} \frac{\eta_2^2}{m_S^2} \left( 1 - \frac{\eta_2^2}{m_S^2} \right)^{3/2} = \frac{\eta$$

This DM still leaves full signatures in the cosmic ray, producing  $2\tau$  or  $4\tau$  which cascade decays into  $e^{\pm}/p^{\pm}/\gamma$ , with good prospects.

In particular, if DM is around 10 GeV, favored by the gamma-ray excesses, the only kinematically accessible mode is  $\tau\tau$ 

$$\langle \sigma v \rangle_{\tau\tau} \approx \frac{\eta_1^2}{4\pi} \left( \frac{\tan\beta \times m_\tau}{m_H^2} \right)^2 \left( 1 - \frac{m_\tau^2}{m_S^2} \right)^{3/2} = 1.2 \times \left( \frac{\eta_1}{0.1} \right)^2 \left( \frac{\tan\beta}{100} \right)^2 \left( \frac{300}{m_H} \right)^4 \, \mathrm{pb.}$$

### Conclusion

- Currently, the null direct detection results place stringent bound on the conventional WIMP
- Indirect detections are also stringent for the light (<100 GeV) DM. But recently, γ-ray excesses are reported, fit by DM DM->ττ
- Quark-phobic dark matter easily avoid LUX etc., and is in accords with the anomaly, and g-2...
- We demonstrate a simple example: 2HDML+S

Thanks for your attention