

Dirac dark matter with a charged mediator:  
a comprehensive one-loop analysis  
of the direct detection phenomenology

Sebastian Wild

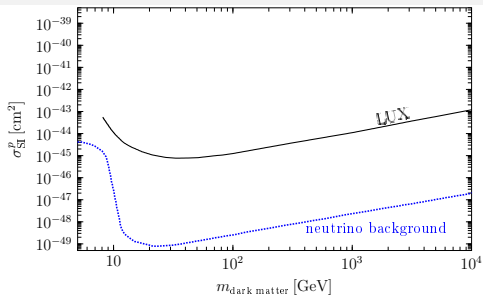
Technische Universität München, Munich



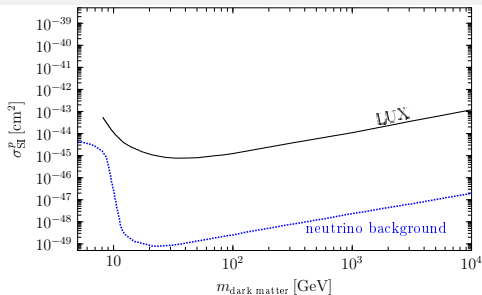
“Exploring the dark sector”  
KIAS, Seoul, March 18, 2015

Based on 1503.03382 (Alejandro Ibarra, SW)

# Loop effects in direct detection: some simple estimates

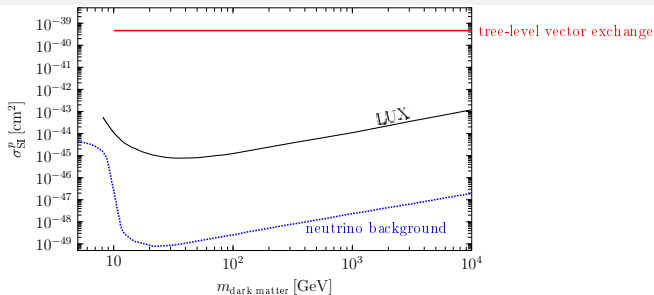


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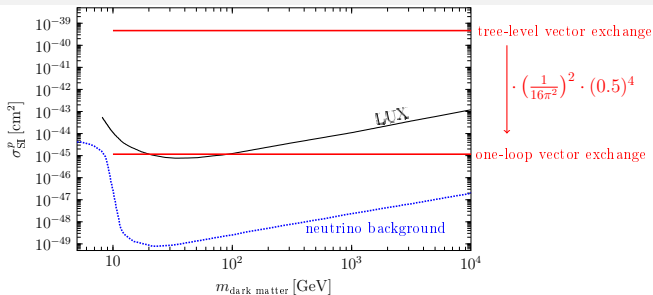
- Consider a **Dirac dark matter** particle  $\chi$
- If  $\chi$  can scatter of nucleons  $N$  at **tree-level** via the exchange of a  $Z$ ...:  
 $\hookrightarrow \mathcal{L}_{\text{eff}} = f_N \bar{\chi} \gamma^\mu \chi \bar{N} \gamma_\mu N$  with  $f_N \sim \frac{G_F}{4\sqrt{2}}$

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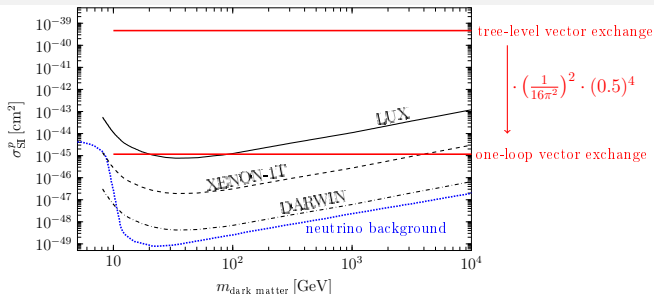
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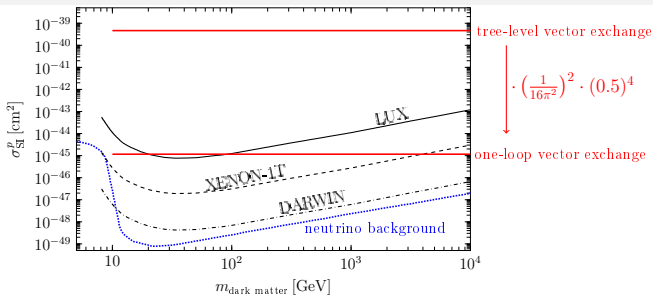
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Direct detection experiments are getting sensitive to loop-induced scattering!

# A simplified model of Dirac dark matter

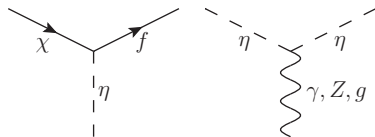
- **Dark matter particle**  $\chi$ : Dirac fermion, SM singlet  $(\mathbf{1}, \mathbf{1})_0$   
 $\hookrightarrow$  no tree-level coupling to  $Z$ ,  $W^\pm$  or  $h$
- **Scalar mediator**  $\eta$   
 $\hookrightarrow$  couples  $\chi$  to one SM fermion  $f_R$  via  $\mathcal{L} = -y\eta^\dagger\bar{\chi}f_R$   
 $\hookrightarrow m_\eta > m_\chi$



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Coupling to...	Quantum numbers of $\eta$
$e_R, \mu_R$ or $\tau_R$	$(\mathbf{1}, \mathbf{1})_{-1}$
$u_R, c_R$ or $t_R$	$(\mathbf{3}, \mathbf{1})_{2/3}$
$d_R, s_R$ or $b_R$	$(\mathbf{3}, \mathbf{1})_{-1/3}$



- Scenarios with coupling to a left-handed SM fermion doublet are very similar  
 $\hookrightarrow$  see the paper for detailed discussion of all the possible cases

# Thermal parameter space

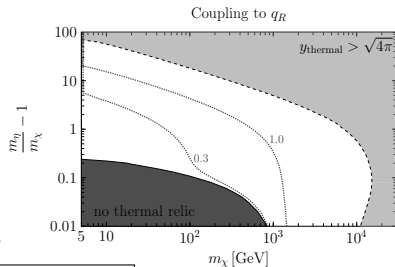
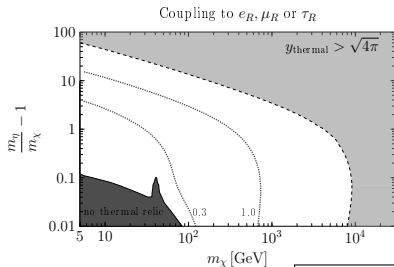
- Only three parameters:  
DM mass  $m_\chi$ , mass splitting  $m_\eta/m_\chi$ , Yukawa coupling  $y$
- The requirement of  $\Omega_\chi h^2 \simeq 0.12$  fixes  $y \equiv y_{\text{thermal}}(m_\chi, m_\eta/m_\chi)$ 
  - ↪ main annihilation channel for the freeze-out:  $\chi\chi \rightarrow f\bar{f}$
  - ↪ For  $m_\eta/m_\chi \lesssim 1.2$ , also coannihilations are important, e.g.  $\eta\eta^\dagger \rightarrow \gamma\gamma$

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For every coupling scenario ( $e_R, u_R, \dots$ ), the **thermal parameter space** is only two-dimensional, spanned by  $m_\chi$  and  $m_\eta/m_\chi$



Highly testable models!

# Intermediate outline

- In the paper, we analyze the direct detection phenomenology for all 15 possible coupling scenarios:

$e_R, (\nu_e, e_L), u_R, d_R, (u_L, d_L) \times$  three generations

- In this talk, I will present the results for coupling to

1)  $u_R$

2)  $\tau_R$

3)  $b_R$

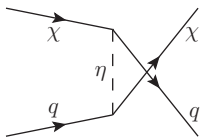
4)  $t_R$

↪ this selection encapsulates most of the interesting features appearing in all possible coupling scenarios

- Some of these cases have (partially) already been discussed:  
Agrawal& [1109.3516,1402.7369,1404.1373], Bai, Berger [1308.0612,1402.6696],  
Chang& [1307.8120,1402.7358], Kopp& [1401.6457]

## Case (1): Coupling to $u_R$

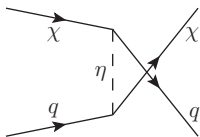
- If  $\chi$  couples to the  $u$  or  $d$  quark, scattering off nuclei is possible at tree-level:



$$\Rightarrow \mathcal{L}_{\text{eff}}^{(\text{SI})} \propto \frac{y^2}{m_\eta^2 - m_\chi^2} \bar{\chi} \gamma^\mu \chi N \gamma_\mu N$$

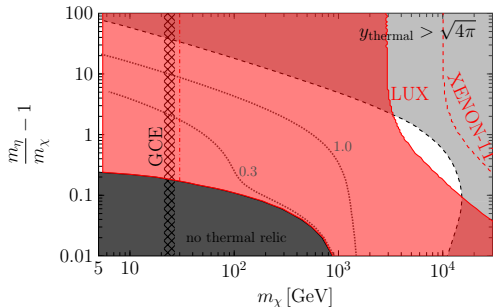
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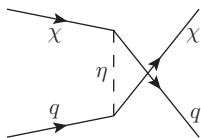
Coupling to  $u_R$



For coupling to  $u_R$  to  $d_R$ ,  
LUX already excludes almost the  
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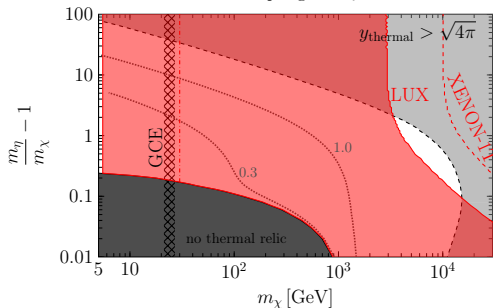
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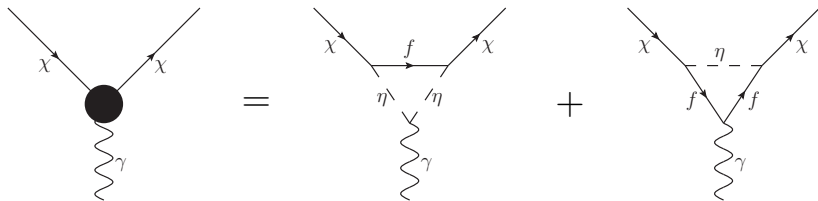
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- However:  $\mathcal{L}_{\text{eff}}^{(\text{SI})} \equiv 0$  for coupling to  $f_R \notin \{u_R, d_R\}$   
 $\hookrightarrow$  in these scenarios, there is **no SI interaction at tree-level!**

## Case (2): Coupling to $\tau_R$ - Leptophilic dark matter

Dark matter coupling to  $\tau_R$  (or any other RH or LH lepton):

- One-loop diagrams induce an **effective DM-photon coupling**:

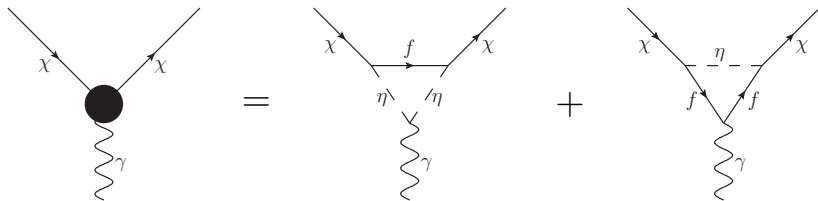




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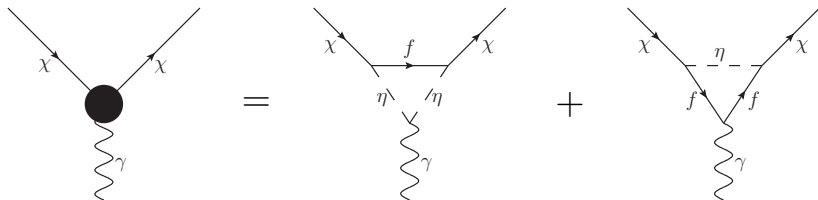
$$\mathcal{L}_{\text{eff}}^{(\gamma)} = \underbrace{\frac{\mu_\chi}{2} \bar{\chi} \sigma^{\mu\nu} \chi F_{\mu\nu}}_{\text{magn. dipole moment}} + \underbrace{b_\chi \bar{\chi} \gamma^\mu \chi \partial^\nu F_{\mu\nu}}_{\text{charge radius}}$$

- $\mu_\chi$  and  $b_\chi$  are calculable (finite) functions of  $m_\chi$ ,  $m_\eta$  and  $y$

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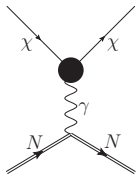
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- $\mu_\chi$  and  $b_\chi$  are calculable (finite) functions of  $m_\chi$ ,  $m_\eta$  and  $y$
- Similar diagrams induce an effective coupling  $\bar{\chi} \chi Z$   
 $\hookrightarrow$  suppressed by  $(m_f/m_\chi)^2 \Rightarrow$  irrelevant for leptophilic dark matter

## Case (2): Coupling to $\tau_R$ - Leptophilic dark matter

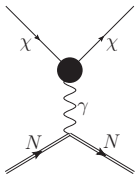


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- Scattering off a nucleus with mass  $m_A$  and spin  $J_A$ :

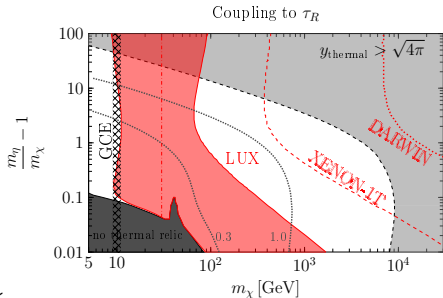
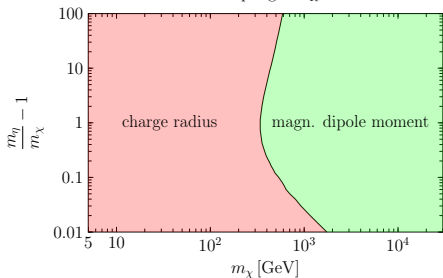
$$\begin{aligned} \frac{d\sigma_{\chi A}}{dE_R} &= \alpha_{\text{em}} \mu_\chi^2 Z^2 \left( \frac{1}{E_R} - \frac{m_A}{2\mu_{\text{red}}^2 v^2} \right) [F_{\text{SI}}(E_R)]^2 && \longrightarrow \text{long-range dipole-charge interaction} \\ &+ \frac{\mu_A^2 \mu_\chi^2 m_A}{\pi v^2} \cdot \frac{J_A + 1}{3J_A} [F_{\text{dipole}}(E_R)]^2 && \longrightarrow \text{dipole-dipole interaction} \\ &&& \propto \mu_A^2 \\ &+ \frac{2m_A Z^2 \alpha_{\text{em}}}{v^2} \left( b_\chi + \frac{\mu_\chi}{2m_\chi} \right)^2 [F_{\text{SI}}(E_R)]^2 && \longrightarrow \text{spin-independent contact interaction} \end{aligned}$$

# Case (2): Coupling to $\tau_R$ - Leptophilic dark matter



Coupling to  $\tau_R$

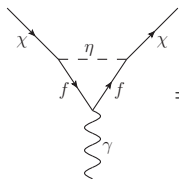
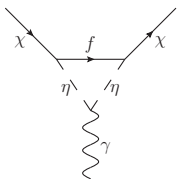
$$\text{via } \mathcal{L}_{\text{eff}}^{(\gamma)} = \underbrace{\frac{\mu_\chi}{2} \bar{\chi} \sigma^{\mu\nu} \chi F_{\mu\nu}}_{\text{magn. dipole moment}} + \underbrace{b_\chi \bar{\chi} \gamma^\mu \chi \partial^\nu F_{\mu\nu}}_{\text{charge radius}}$$



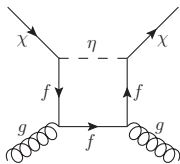
- LUX already excludes a thermal relic for  $m_\chi \lesssim 100$  GeV
- Future DD experiments can completely close in on that scenario
  - ↪ ... even for  $m_\chi \gtrsim 1$  TeV
  - ↪ ... despite the fact that the scattering is only loop-induced!

### Case (3): Coupling to $b_R$

- Again, there is no tree-level scattering off nuclei
- Relevant one-loop diagrams:



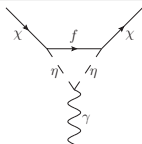
$\Rightarrow$  charge radius & magnetic dipole moment of  $\chi$



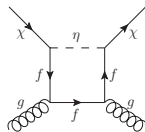
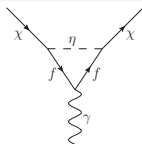
$\Rightarrow$  effective dark matter-gluon coupling  $\bar{\chi}\chi G^{\mu\nu} G_{\mu\nu}$

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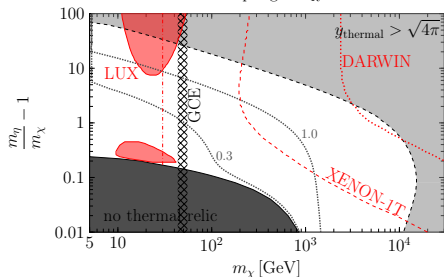
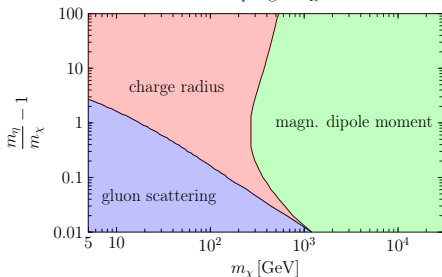
## Case (3): Coupling to $b_R$



Coupling to  $b_R$

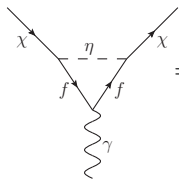
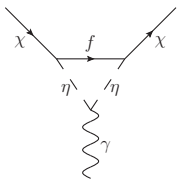


Coupling to  $b_R$

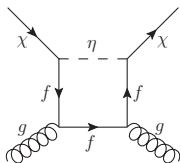


- Excellent prospects for future direct detection experiments
- In this model, the dominant annihilation mode is  $\chi\bar{\chi} \rightarrow b\bar{b}$ 
  - ↪ for  $m_{DM} \simeq 50$  GeV, this model could explain the **Galactic Center Excess** observed by Fermi
  - ↪ well within the reach of XENON-1T!

## Case (4): Coupling to $t_R$

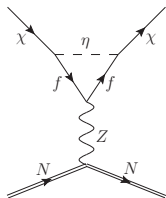
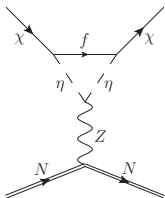


$\Rightarrow$  charge radius & magnetic dipole moment of  $\chi$



$\Rightarrow$  effective dark matter-gluon coupling  $\bar{\chi}\chi G^{\mu\nu} G_{\mu\nu}$

$$\hookrightarrow \mathcal{L}_{\text{eff}} = f_{S, \text{gluon}} \bar{\chi}\chi \bar{N}N$$

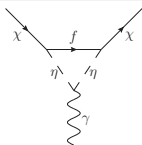


$\Rightarrow$  effective  $\bar{\chi}\chi Z$  vertex

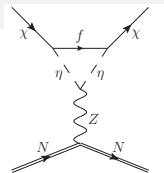
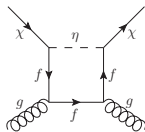
$$\hookrightarrow \mathcal{L}_{\text{eff}} = f_{V,Z} \bar{\chi}\gamma^\mu \chi \bar{N}\gamma_\mu N \text{ with}$$

$$f_{V,Z} \propto (m_f/m_\chi)^2$$

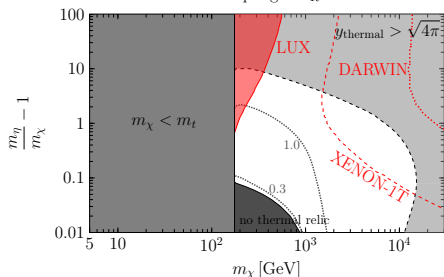
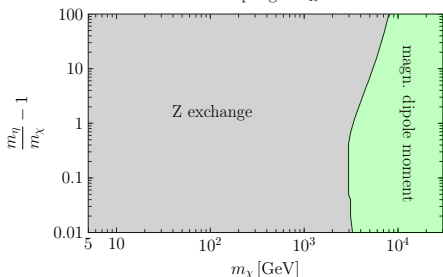
## Case (4): Coupling to $t_R$



Coupling to  $t_R$



Coupling to  $t_R$

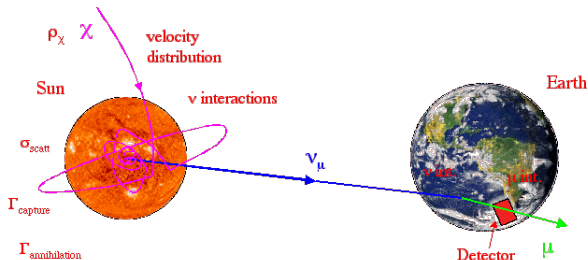


- For DM coupling to the top-quark, the one-loop  $Z$  exchange is dominant for  $m_\chi \lesssim 5$  TeV
- LUX only constrains a rather small part of the thermal parameter space
- However: XENON-1T will be sensitive to dark matter masses  $\gtrsim 1$  TeV



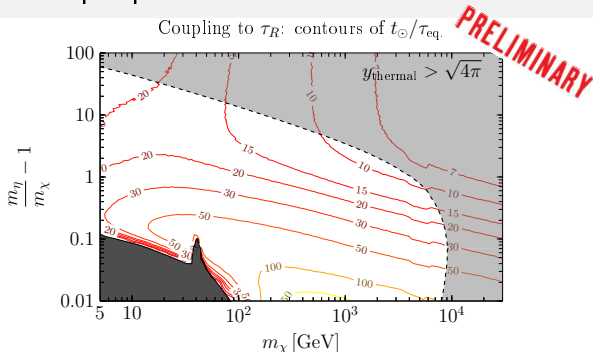
# Outlook: leptophilic dark matter in the Sun?

Can dark matter be captured and annihilate in the Sun,  
if it is **purely leptophilic**?



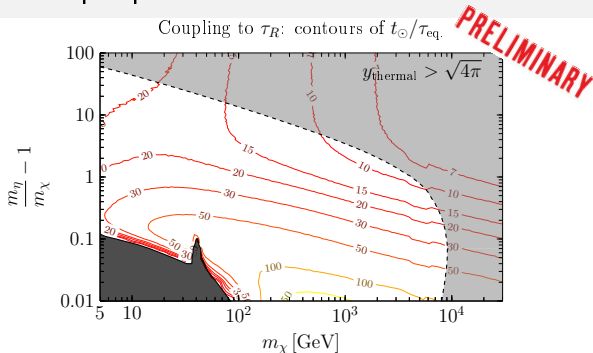
- The IceCube/Super-K limits can be very strong for leptophilic annihilations, e.g.  $\text{DM DM} \rightarrow \tau^+ \tau^-$  or  $\text{DM DM} \rightarrow \nu \bar{\nu}$
- The magnetic dipole moment of  $\chi$  leads to a dipole-dipole interaction with protons  
↳ similar to spin-dependent scattering off protons

# Equilibration of leptophilic Dark Matter in the Sun



- Again, we fix the coupling  $y$  to its thermal value ( $\Omega_{\text{DM}} h^2 = 0.12$ )
- $A_\odot = \frac{1}{2} C_\odot \tanh^2(t_\odot/\tau_{\text{eq}})$  with  $\tau_{\text{eq}} \propto \frac{1}{\sqrt{\sigma_{\text{capture}} \cdot \langle \sigma v \rangle}}$ 
  - ↪ **Equilibrium condition:**  $t_\odot/\tau_{\text{eq}} \gg 1$
  - ↪ contours show constant values of  $t_\odot/\tau_{\text{eq}}$

# Equilibration of leptophilic Dark Matter in the Sun

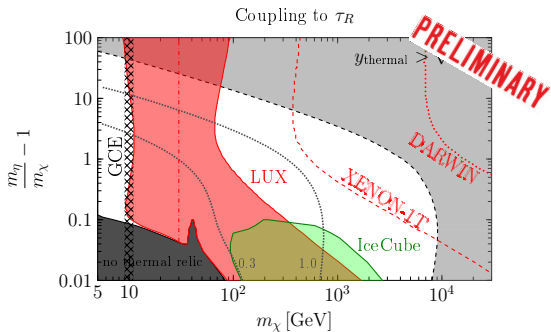


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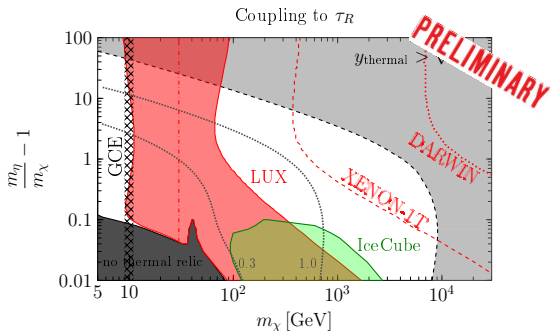


In our simplified model of leptophilic Dirac Dark Matter, we indeed have equilibration in the whole thermal parameter space

# IceCube and LUX limits on leptophilic Dirac dark matter



# IceCube and LUX limits on leptophilic Dirac dark matter



## Complementarity between LUX and IceCube:

- ↪ For large DM masses and small mass splitting of  $\chi$  and  $\eta$ , IceCube is stronger than LUX
- ↪ LUX and IceCube are affected differently by systematics, e.g. the dark matter velocity distribution

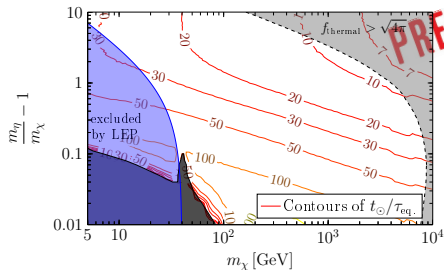
# Conclusions

- Direct detection experiments have already reached the necessary sensitivity to probe dark matter-nucleon interactions mediated **only at one-loop order**
- We demonstrated this in the context of a simplified model with a singlet Dirac fermion as the dark matter particle
  - ↔ for **any** possible coupling scenario of dark matter, LUX already excludes part of the thermal parameter space
  - ↔ future direct detection experiments have excellent prospects to probe **large parts** of the parameter space, even multi-TeV dark matter
  - ↔ this approach is complementary and competitive to other searches:
    - LHC: monojet, direct production of  $\eta$ , ...
    - indirect detection with  $\gamma$ 's from dwarf galaxies
    - for leptophilic DM: indirect detection with  $e^+/e^-$  (AMS-02)
- Furthermore, via the one-loop scattering diagrams, leptophilic dark matter can be captured in the Sun
  - ↔ limits from IceCube can be competitive with LUX in some parts of the parameter space

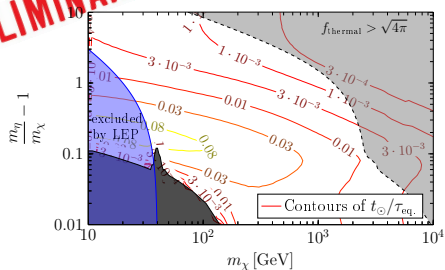
Backup slides

# Equilibration of leptophilic Dark Matter in the Sun

Dirac DM



Majorana DM

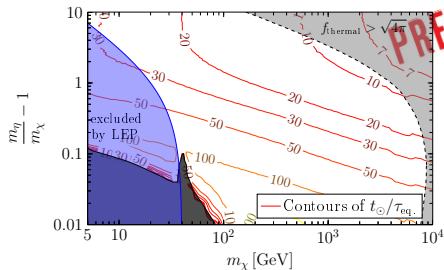


- For this plot, we fix the coupling  $f$  to its **thermal value** ( $\Omega_{\text{DM}} h^2 = 0.12$ )

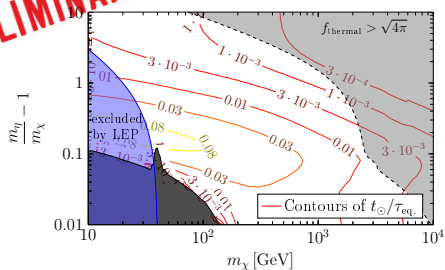


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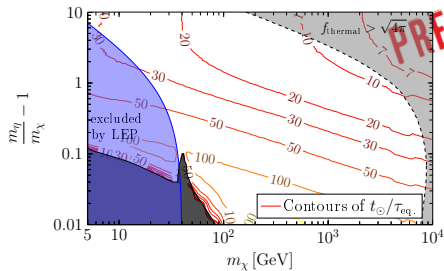
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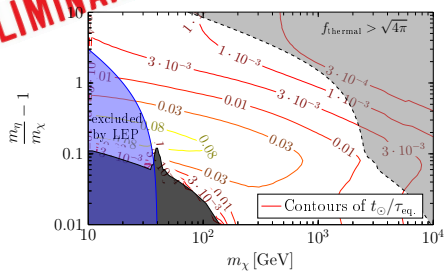
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In our simplified model of leptophilic Dark Matter...

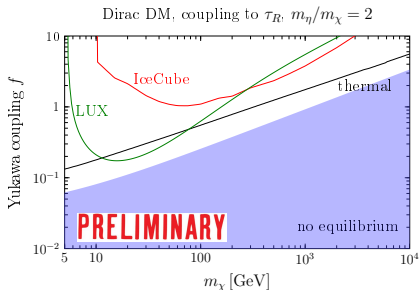
- DM **is** in equilibrium in the Sun if it is a Dirac fermion
- DM **is not** in equilibrium in the Sun if it is a Majorana fermion

# IceCube and LUX limits on leptophilic Dirac Dark Matter

- Dirac Dark Matter annihilates into  $f\bar{f}$  (no helicity suppression)
  - ↪ **strong neutrino signal** for DM coupling to  $\tau_R$ ,  
or to any left-handed lepton doublet
  - ↪ for coupling to  $e_R$  or  $\mu_R$  neutrinos (only) from weak FSR

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Leptophilic Dirac DM coupling to  $\tau_R$  or to any left-handed lepton doublet:  
For  $m_\chi \gtrsim 100$  GeV, IceCube limits can be **stronger than LUX**

- Black line: Yukawa coupling  $f$  leading to  $\Omega h^2 = 0.12$
- Remark: Leptophilic Majorana Dark Matter is not in equilibrium in the Sun
  - ↪ no meaningful limits from IceCube