

Plan of Lecture

- Part 0:
General introduction
- Part 1:
Fluctuation theorems
- **Part 2:**
Classical/quantum measurement and information
- Part 3:
Information thermodynamics

Part 2:

Classical/quantum measurement and information

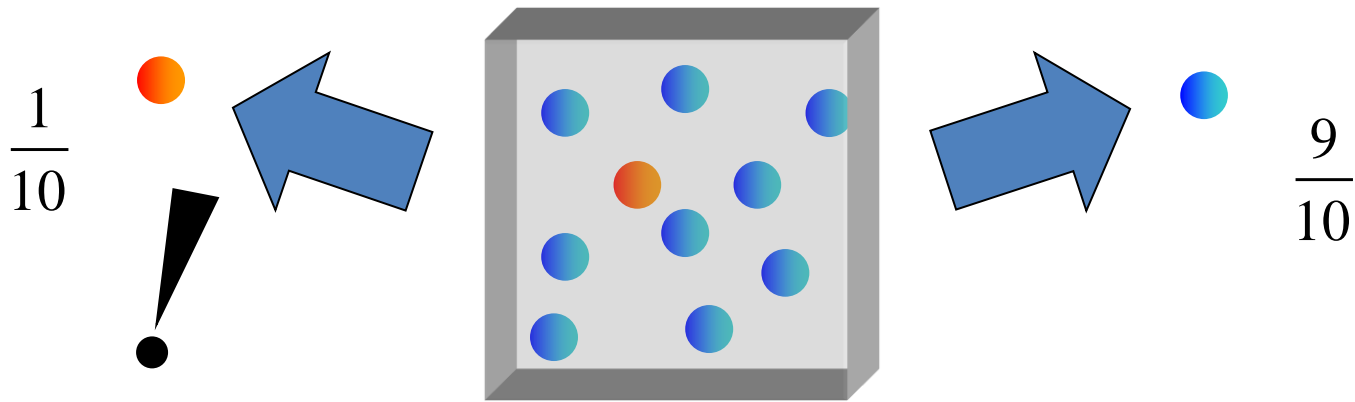
- Shannon information
- Mutual information
- Quantum measurement
- Quantum mutual information
- QC-mutual information

Part 2:

Classical/quantum measurement and information

- **Shannon information**
- Mutual information
- Quantum measurement
- Quantum mutual information
- QC-mutual information

Shannon Information (1)



Information content with event k : $\ln \frac{1}{p_k}$

Average 

Shannon information: $H = \sum_k p_k \ln \frac{1}{p_k}$

Shannon Information (2)

$$0 \leq H \leq \ln N$$

N : the number of k 's

$$\begin{aligned} p_i &= 1 \quad \text{for a single } i \\ p_k &= 0 \quad \text{for } k \neq i \end{aligned}$$

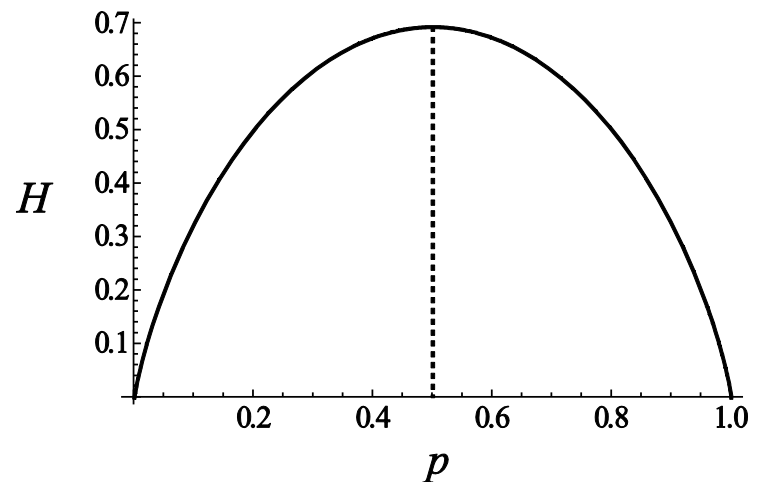
$$\begin{aligned} p_k &= 1/N \\ &\text{for all } k \end{aligned}$$

$$H = -\sum_k p_k \ln p_k$$

Ex. Binary system "0": probability p
 "1": probability $1-p$

$$H = -p \ln p - (1-p) \ln(1-p)$$

$$0 \leq H \leq \ln 2$$



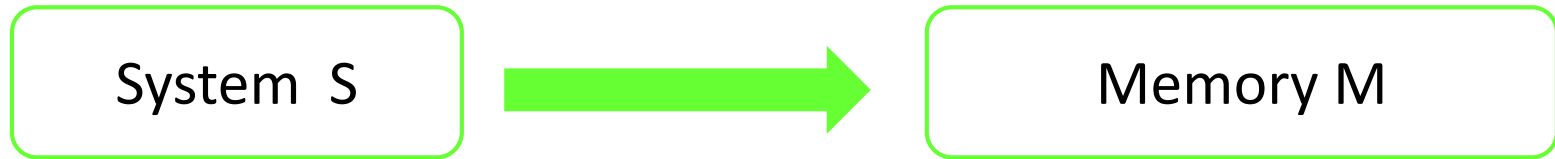
Characterizes the **randomness** of the system

Part 2:

Classical/quantum measurement and information

- Shannon information
- **Mutual information**
- Quantum measurement
- Quantum mutual information
- QC-mutual information

Mutual Information (1)



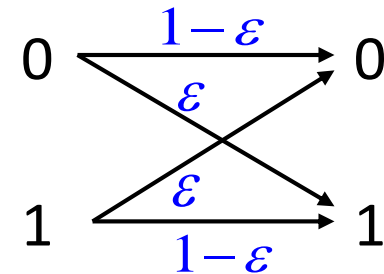
Measurement or communication
(with stochastic error, in general)

$p(s)$: distribution of the measured state of S

$p(m)$: distribution of the outcome in M

$p(m | s)$: conditional probability characterizing the error

$p(s, m) = p(m | s)p(s)$: joint distribution of S and M

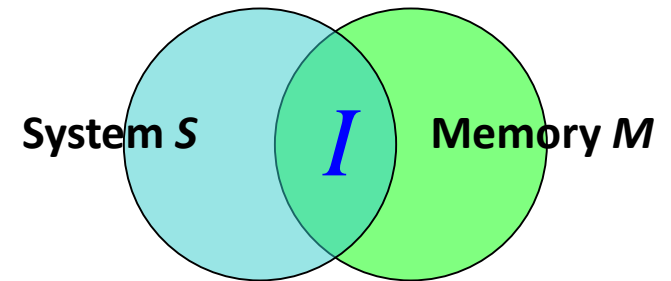


Ex. Binary symmetric channel

Mutual Information (2)

$$I(S : M) \equiv H(S) + H(M) - H(SM)$$

$$= \sum_{sm} p(s, m) \ln \frac{p(s, m)}{p(s) p(m)}$$



$p(s)$: distribution of the measured state of S

$p(m)$: distribution of the outcome in M

$p(m | s)$: conditional probability characterizing the error

$p(s, m) = p(m | s) p(s)$: joint distribution of S and M

$$H(S) = -\sum_s p(s) \ln p(s) \quad H(M) = -\sum_m p(m) \ln p(m) \quad H(SM) = -\sum_{sm} p(s, m) \ln p(s, m)$$

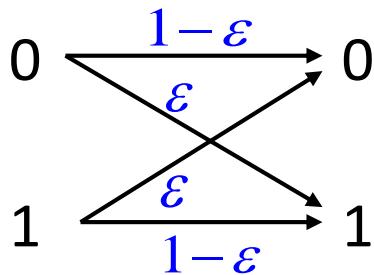
Mutual Information (3)

$$I(S : M) \equiv H(S) + H(M) - H(SM)$$

$$0 \leq I \leq H(M)$$

No information

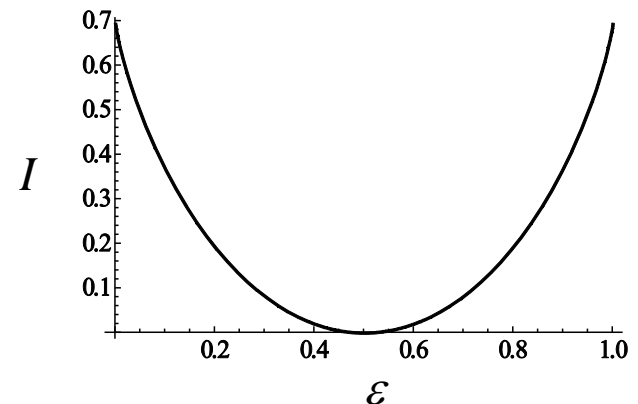
No error



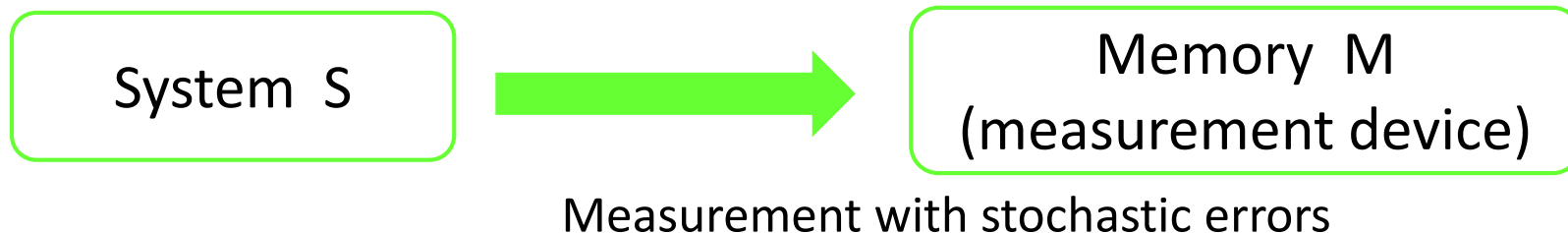
Ex. Binary symmetric channel

Correlation between S and M

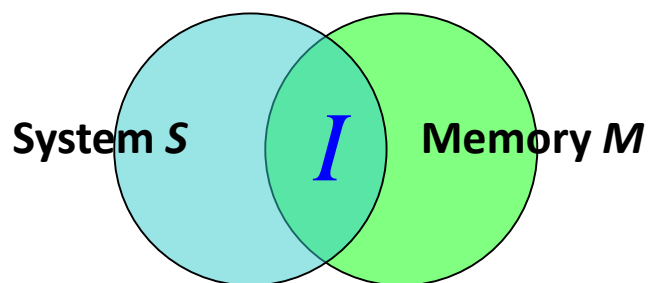
$$I = \ln 2 + \epsilon \ln \epsilon + (1 - \epsilon) \ln(1 - \epsilon)$$



Mutual Information: Summary



$$I(S : M) \equiv H(S) + H(M) - H(SM)$$



$$0 \leq I \leq H(M)$$

No information

Error-free

Shannon information:
Randomness of the system

Mutual information:
Correlation between the system
and the memory

Part 2:

Classical/quantum measurement and information

- Shannon information
- Mutual information
- **Quantum measurement**
- Quantum mutual information
- QC-mutual information

Quantum State

Density operator: $\rho = \sum_i q_i |\varphi_i\rangle\langle\varphi_i|$

Statistical mixture of state vector $|\varphi_i\rangle$ with probability q_i

Non-negativity: $\rho \geq 0$

Unity: $\text{tr}[\rho] = \sum_i q_i = 1$

Quantum Measurement

Projection measurement (error-free)

Observable $A = \sum_k \alpha_k P_k$

Measurement outcome k

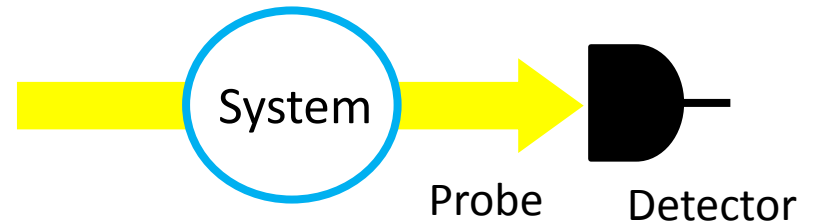
Projection operators $\{P_k\}$

$$\sum_k P_k = I$$

Probability $p_k = \text{tr}[\rho P_k]$

Born's rule

Indirect measurement (with error)



Measurement outcome k

POVM $\{E_k\}$ (positive operators)

$$\sum_k E_k = I$$

Probability $p_k = \text{tr}[\rho E_k]$

Projection Measurement

Projection measurement on density operator $\rho = \sum_i q_i |\varphi_i\rangle\langle\varphi_i|$

Observable $A = \sum_k \alpha_k P_k$ P_k : projection operator

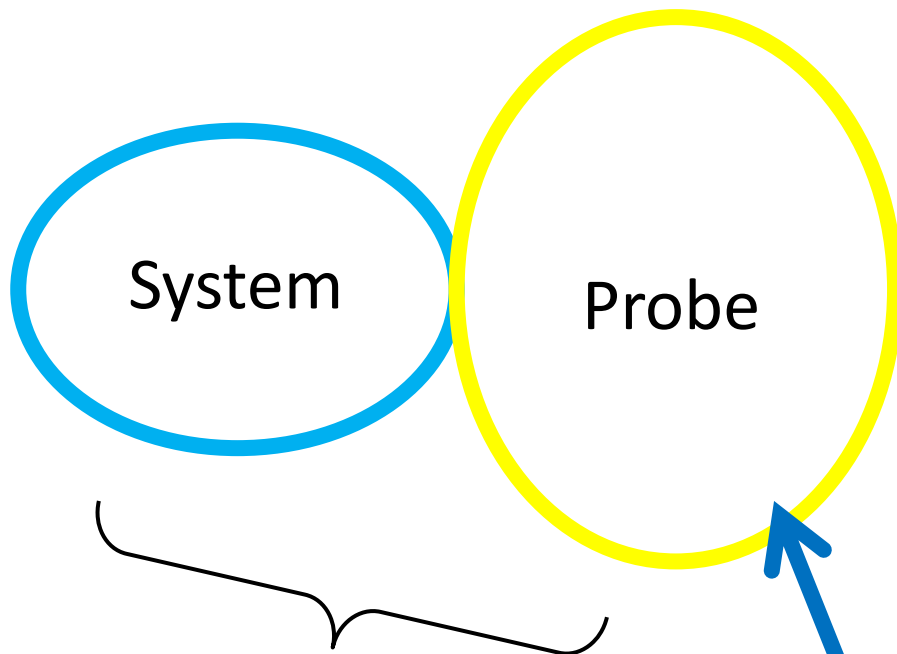
Probability of outcome k

Born's rule $p_k = \sum_i q_i \langle\varphi_i|P_k|\varphi_i\rangle = \text{tr}[P_k\rho]$

Post-measurement state with outcome k

Projection postulate $\frac{1}{p_k} P_k \rho P_k$

Schematic of Indirect Measurement




Unitary interaction of the system and the probe

Projection measurement on the probe

$$P_k = |\psi_k\rangle\langle\psi_k| \quad \{\psi_k\} : \text{orthonormal basis}$$

$$\rho \otimes \rho_P$$



$$U \rho \otimes \rho_P U^\dagger$$

Assume: $\rho_P = |\psi\rangle\langle\psi|$

Probability of Indirect Measurement

Probability of outcome k :

$$\begin{aligned} p_k &= \text{tr}_{\text{SP}} [P_k U \rho \otimes |\psi\rangle\langle\psi| U^\dagger P_k] \\ &= \text{tr}_S [\langle\psi_k| U \rho \otimes |\psi\rangle\langle\psi| U^\dagger |\psi_k\rangle] \\ &= \text{tr}_S [\langle\psi_k| U |\psi\rangle \rho \langle\psi| U^\dagger |\psi_k\rangle] \\ &= \text{tr}_S [M_k \rho M_k^\dagger] \\ &= \text{tr}_S [E_k \rho] \end{aligned}$$

POVM:

$$E_k = M_k^\dagger M_k$$

Kraus operator:

$$M_k \equiv \langle\psi_k| U |\psi\rangle$$

$$\sum_k p_k = 1 \quad \longleftrightarrow \quad \sum_k E_k = I_d$$

Back-action of Indirect Measurement

Post-measurement state with outcome k :

$$\begin{aligned}\rho_k &= \frac{1}{p_k} \text{tr}_P [P_k U \rho \otimes |\psi\rangle\langle\psi| U^\dagger P_k] \\ &= \frac{1}{p_k} \langle \psi_k | U \rho \otimes |\psi\rangle\langle\psi| U^\dagger | \psi_k \rangle \\ &= \frac{1}{p_k} M_k \rho M_k^\dagger\end{aligned}$$
$$M_k \equiv \langle \psi_k | U | \psi \rangle$$

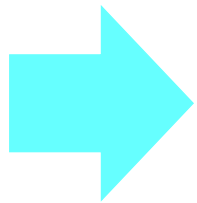
Example 1: Projection Measurement

Two-level atom

Kraus operators: $M_0 = |0\rangle\langle 0|$ $M_1 = |1\rangle\langle 1|$ (projection)

POVM: $E_0 = |0\rangle\langle 0|$ $E_1 = |1\rangle\langle 1|$ (projection)

$$\rho = q_0 |0\rangle\langle 0| + q_1 |1\rangle\langle 1| + \alpha |0\rangle\langle 1| + \alpha^* |1\rangle\langle 0|$$

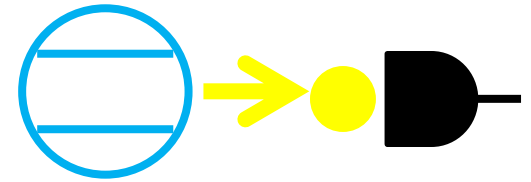


$$p_0 = q_0 \quad p_1 = q_1 \quad (\text{error-free})$$

$$\rho_0 = |0\rangle\langle 0| \quad \rho_1 = |1\rangle\langle 1| \quad (\text{projection})$$

Example 2: Measurement by Photon Emission

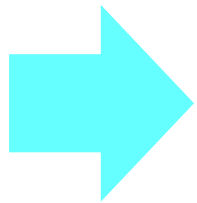
Two-level atom



Kraus operators: $M_0 = |0\rangle\langle 0|$ $M_1 = |0\rangle\langle 1|$ (not projection)

POVM: $E_0 = |0\rangle\langle 0|$ $E_1 = |1\rangle\langle 1|$ (projection)

$$\rho = q_0 |0\rangle\langle 0| + q_1 |1\rangle\langle 1| + \alpha |0\rangle\langle 1| + \alpha^* |1\rangle\langle 0|$$



$$p_0 = q_0 \quad p_1 = q_1 \quad (\text{error-free})$$

$$\rho_0 = |0\rangle\langle 0| \quad \rho_1 = |0\rangle\langle 0| \quad (\text{not projection})$$

Example 3: Measurement with Error

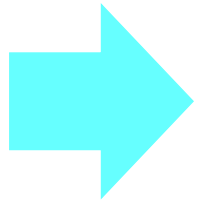
Two-level atom

Kraus operators: $M_0 = \sqrt{1-\varepsilon}|0\rangle\langle 0| + \sqrt{\varepsilon}|1\rangle\langle 1|$ $M_1 = \sqrt{\varepsilon}|0\rangle\langle 0| + \sqrt{1-\varepsilon}|1\rangle\langle 1|$

POVM: $E_0 = (1-\varepsilon)|0\rangle\langle 0| + \varepsilon|1\rangle\langle 1|$ $E_1 = \varepsilon|0\rangle\langle 0| + (1-\varepsilon)|1\rangle\langle 1|$

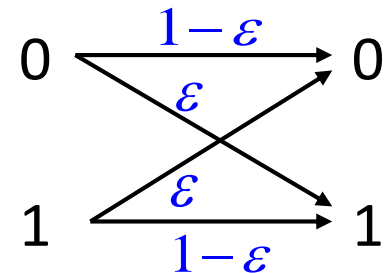
ε : error rate

$$\rho = q_0|0\rangle\langle 0| + q_1|1\rangle\langle 1| + \alpha|0\rangle\langle 1| + \alpha^*|1\rangle\langle 0|$$



$$p_0 = (1-\varepsilon)q_0 + \varepsilon q_1$$

$$p_1 = \varepsilon q_0 + (1-\varepsilon)q_1$$

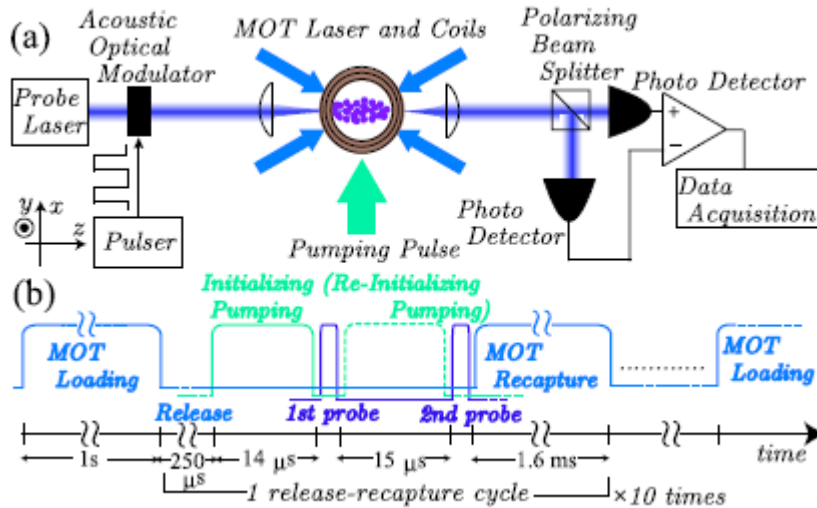


Real Example: QND Measurement of Atomic Spin

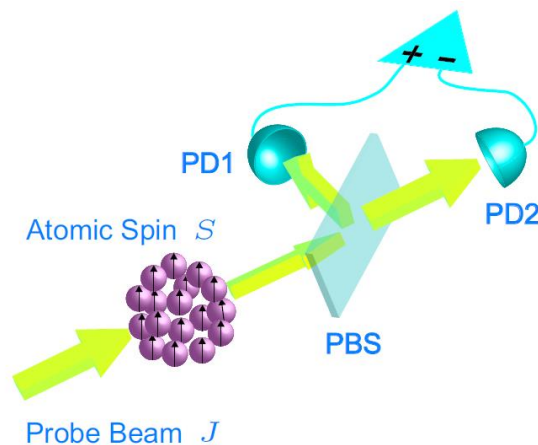
Takano *et al.*, Phys. Rev. Lett. **102**, 033601 (2009)

QND = Quantum Non-Demolition

System = Atomic Nuclear Spins
 Probe = Polarized photons

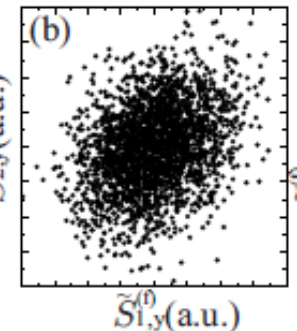
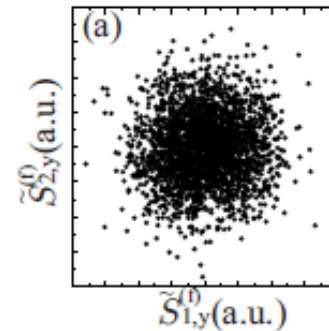


Correlation of two-time measurements



Without atoms

With atoms



No correlation

Correlation by wave function collapse

Quantum Measurement: Summary

Projection measurement (error-free)

Observable $A = \sum_k \alpha_k P_k$

Measurement outcome k

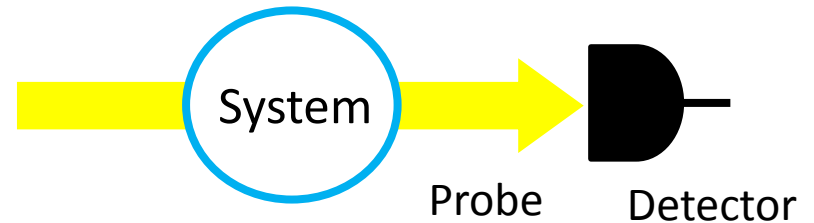
Projection operators $\{P_k\}$

$$\sum_k P_k = I$$

Probability $p_k = \text{tr}[\rho P_k]$

Born's rule

Indirect measurement (with error)



Measurement outcome k

POVM $\{E_k\}$ (positive operators)

$$\sum_k E_k = I$$

Probability $p_k = \text{tr}[\rho E_k]$

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Classical/quantum measurement and information

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Quantum Entropy

Von Neumann entropy: $S(\rho) = -\text{tr}[\rho \ln \rho]$

ρ : density operator

$$\rho = \sum_i q_i |\varphi_i\rangle\langle\varphi_i| \quad \rightarrow \quad S(\rho) = -\sum_i q_i \ln q_i$$

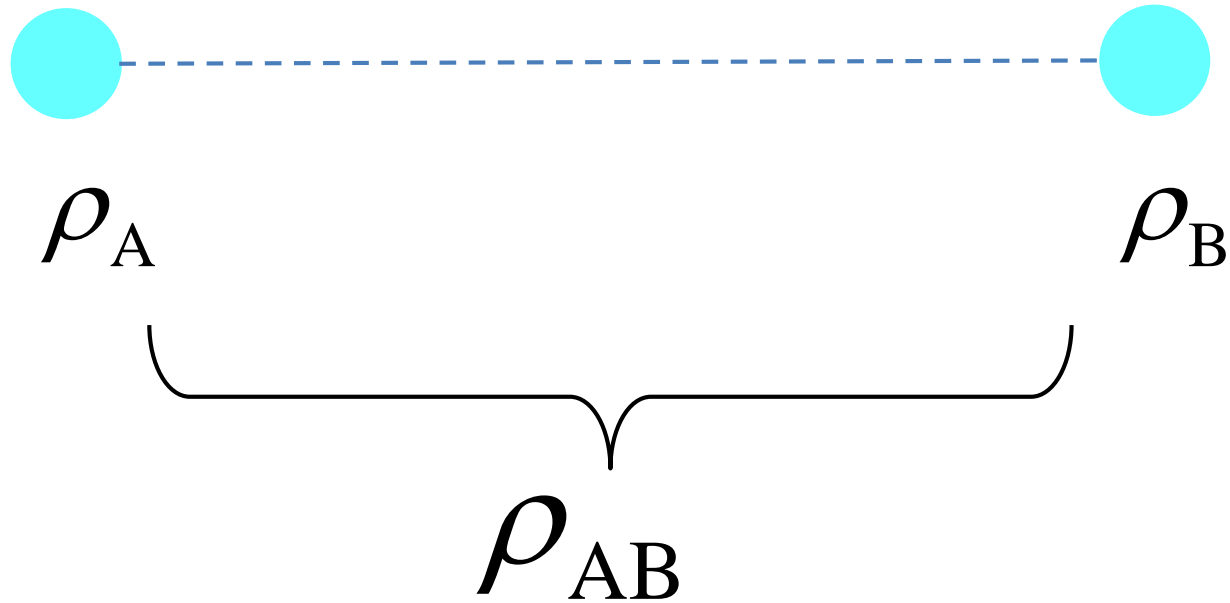
with an orthonormal basis

**Characterizes the randomness of the classical mixture
in the density operator**

Variety of Quantum Information

- Quantum mutual information
- Quantum discord
- Holevo's chi
- QC-mutual information (Groenewold information)

Bipartite Quantum System



$$\rho_A = \text{tr}_B[\rho_{AB}] \quad \rho_B = \text{tr}_A[\rho_{AB}]$$

Quantum Mutual Information (1)

$$I_{A:B}(\rho_{AB}) = S(\rho_A) + S(\rho_B) - S(\rho_{AB})$$

$$I_{A:B}(\rho_{AB}) \geq 0 \quad \longleftrightarrow \quad S(\rho_A) + S(\rho_B) \geq S(\rho_{AB})$$

(Subadditivity of von Neumann entropy)

Purely classical correlation \Rightarrow Classical mutual information

$$\rho_{AB} = \sum_{xy} q_{xy} |\varphi_x\rangle\langle\varphi_x| \otimes |\phi_y\rangle\langle\phi_y| \quad \text{with orthonormal bases}$$

$$\longrightarrow \quad I_{A:B}(\rho_{AB}) = \sum_{xy} q_{xy} \ln \frac{q_{xy}}{q_x q_y}$$

Quantum Mutual Information (2)

Example: 2 qubits

$$\rho_{AB} = |00\rangle\langle 00| \text{ (no correlation)} \quad \rightarrow \quad I_{A:B}(\rho_{AB}) = 0$$

$$\rho_{AB} = \frac{1}{2} (|00\rangle\langle 00| + |11\rangle\langle 11|) \quad \rightarrow \quad I_{A:B}(\rho_{AB}) = \ln 2$$

(classical correlation)

$$\rho_{AB} = \frac{1}{2} (|00\rangle\langle 00| + |11\rangle\langle 11| + |11\rangle\langle 00| + |00\rangle\langle 11|)$$

(entanglement)

$$\rightarrow I_{A:B}(\rho_{AB}) = 2 \ln 2$$

Part 2:

Classical/quantum measurement and information

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QC-mutual Information (1)

Information flow from **Q**uantum system to **C**lassical outcome by quantum measurement

$$I_{\text{QC}} \equiv S(\rho) - \sum_y p_y S(\rho_y)$$

ρ : measured density operator

$p(y) = \text{tr}[\rho M_y^\dagger M_y]$: probability of obtaining outcome k

$\rho_y = \frac{1}{p(y)} M_y \rho M_y^\dagger$: post-measurement state with outcome k

Assumed a single Kraus operator for each outcome

H. J. Groenewold, Int. J. Theor. Phys. **4**, 327 (1971).

M. Ozawa, J. Math. Phys. **27**, 759 (1986).

TS and M. Ueda, PRL **100**, 080403 (2008).

QC-mutual Information (2)

$$0 \leq I_{\text{QC}} \leq H$$

$$H = -\sum_y p(y) \ln p(y)$$

No information

Any POVM
element is
identity operator

Error-free & classical

Any POVM element is
projection and commutable
with measured state

Classical measurement, I_{QC} reduces to the classical mutual information

If the measured state is a pure state: $I_{\text{QC}} = 0$

QC-mutual Information (3)

The QC-mutual information gives an upper bound of the accessible classical information x encoded in ρ

F. Buscemi, M. Hayashi, and M. Horodecki, PRL **100**, 210504 (2008).

$$\rho = \sum_x q(x) \rho_x \quad : \text{an arbitrary decomposition}$$

ρ_x are not necessarily orthogonal

$$p(x, y) = \text{tr}[M_y^\dagger M_y \rho_x] q(x) \quad p(y) = \text{tr}[M_y^\dagger M_y \rho] = \sum_x p(x, y)$$

$$I = \sum_{xy} p(x, y) \ln \frac{p(x, y)}{q(x) p(y)}$$

Theorem: $I \leq I_{\text{QC}}$

A variant (a “dual”) of the Holevo bound