Current status of $a_{\mu}^{\rm SM}$

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Two-loop corrections to $(g - 2)_{\mu}$ in SM und MSSM

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Motivation

$$a_{\mu}^{ extsf{E821}} = (11659208.9 \pm 6.3) imes 10^{(-10)}$$
 [Bennett et al. '06]

New experiment: moving from BNL to Fermilab E969 $0.54 \rightarrow 0.14 \text{ ppm}$

current accuracy of $a^{SM}_{\mu}:~0.42~\text{ppm}$ $_{\text{[Davier et al. '10]}}$

motivated to improve the accuracy of all aspects of the theory prediction, in SM and SUSY



Summary

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Current status of $a_\mu^{ m SM}$

MSSM two-loop corrections

Numeric and leading logarithmic approximation

Summary



Recent SM theory progress

- QED 5-loop calculation completed [Aoyama, Hayakawa, Kinoshita, Nio '12]
- Convergence of hadronic contributions [Davier et al., Hagiwara et al., Benayoun et al.]
- Electroweak (full 2-loop) contributions with $M_H = 126 \text{GeV}$:

 $a_{\mu}^{EW} = (153.6 \pm 1.0) imes 10^{-11}$ [Gnendiger, Stöckinger, S-K '13]



approx.[Czarnecki, Krause, Marciano]

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a_{μ} as important constraint on SUSY

$$\Delta a_{\mu}({
m E821-SM}) = egin{cases} (28.7\pm8.0) imes10^{-10} \ {}_{ ext{[Davier et al.]}}\ (26.1\pm8.0) imes10^{-10} \ {}_{ ext{[Hagiwara et al.]}} \end{cases}$$

 $\sim 3\sigma$ deviation

tension between LHC-limits and a_{μ}

- a_{μ} is an important constraint on SUSY
- some scenarios (e.g. Constrained MSSM) cannot explain a_{μ} deviation
- still, large contributions possible, e.g. if sleptons << squarks (non-traditional models)

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Already known SUSY corrections



• SUSY one-loop diagram: μ , M_1 , M_2 , M_E , M_L , $\tan \beta$ [Fayet '80]...[Moroi '96]

• SUSY two-loop corrections to SM 1L diagrams

[Chen, Geng '01], [Arhib, Baek '02], [Heinemeyer, Stöckinger, Weiglein '03, '04]

• Photonic corrections to SUSY 1L diagrams: μ , M_1 , M_2 , M_E , M_L , tan β

[v. Weitershausen, Schäfer, Stöckinger, S-K '10]

estimated resulting MSSM theory error $\sim 3 \times 10^{-10}$ $_{\rm [Stöckinger '06]}$

MSSM photonic two-loop corrections



[v. Weitershausen, Schäfer, Stöckinger, S-K '10]

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- exact results of photonic two-loop corrections
- same parameter dependence as SUSY one-loop diagram
- reproduces the leading QED-logarithms: $\frac{4\alpha}{\pi} \log \frac{m_{\mu}}{M_{SUSY}} a_{\mu}^{1LSUSY}$ [Degrassi, Giudice '98]
- - (7...9)% corrections for $100 < M_{SUSY} < 1000 \,{
 m GeV}$

MSSM fermion/sfermion two-loop corrections

[Fargnoli, Gnendiger, Paßehr, Stöckinger, S-K '13]

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- remaining class with dependence on squarks
- $(f, \tilde{f}) \in \{(\nu, \tilde{\nu}), (I, \tilde{I}), (u, \tilde{u}), (d, \tilde{d})\}, (f', \tilde{f}) \in \{(I, \tilde{\nu}), (\nu^c, \tilde{I}^{\dagger}), (d, \tilde{u}), (u^c, \tilde{d}^{\dagger})\}$
- maximum complexity: 5 heavy + 2 light scales
- computed exactly, including renormalization
- additional parameter dependence: $M_{Q_i}, M_{U_i}, M_{D_i}, M_{L_i}, M_{E_i}, \quad i \in \{1, 2, 3\}$

MSSM fermion/sfermion two-loop corrections





- iterated one-loop calculation
- 5-mass-scale two-loop diagrams and their counterterm diagrams
- the sum of the inner loops build up compact vertices: $i\Gamma_{ji}^{0/-}$, $i\Sigma_{ii}$

MSSM fermion/sfermion two-loop corrections: Γ_{ji}^{0}

[arXiv:1311.1775][Fargnoli, Gnendiger, Paßehr, Stöckinger, S-K '13]



$$\begin{split} \Gamma^{0\mu}_{ij\tilde{f}_{k}}(\ell) &= \frac{1}{16\pi^{2}} e Q_{f} \int_{0}^{1} \frac{\mathrm{d} w}{2} \left[\left(\overline{\mathcal{A}}_{ij\tilde{f}_{k}}^{n+} - \overline{\mathcal{A}}_{ij\tilde{f}_{k}}^{n-} \gamma^{5} \right) \frac{\ell \not g \gamma^{\mu} - \ell q^{\mu} + \not g \ell^{\mu} - (\ell \cdot q) \gamma^{\mu}}{\mathcal{D}_{f\tilde{f}_{k}}(\ell)} \\ &+ \left(\overline{\mathcal{B}}_{ij\tilde{f}_{k}}^{n+} - \overline{\mathcal{B}}_{ij\tilde{f}_{k}}^{n-} \gamma^{5} \right) \frac{m_{f}}{w} \frac{\not g \gamma^{\mu} - q^{\mu}}{\mathcal{D}_{f\tilde{f}_{k}}(\ell)} \right]. \end{split}$$



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MSSM fermion/sfermion two-loop corrections: sample analytic result

[arXiv:1311.1775][Fargnoli, Gnendiger, Paßehr, Stöckinger, S-K '13]

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$$\begin{split} \mathbf{a}_{\mu \ ijm\tilde{t}_{k}}^{(nv)} &= \int_{0}^{1} \mathrm{d}\mathbf{w} \left[\mathcal{A}_{jj\tilde{t}_{k}}^{n+} \left(\overline{\mathcal{A}}_{ij\tilde{t}_{k}}^{n+} \mathcal{T}_{AA}^{nv+} + \overline{\mathcal{B}}_{ij\tilde{t}_{k}}^{n+} \mathcal{T}_{AB}^{nv+} \right) + \mathcal{B}_{jj\tilde{t}_{k}}^{n+} \mathcal{T}_{BA}^{nv+} + \overline{\mathcal{B}}_{ij\tilde{t}_{k}}^{n+} \mathcal{T}_{BB}^{n+} \right) \\ &+ \mathcal{A}_{jj\tilde{\mu}m}^{n-} \left(\overline{\mathcal{A}}_{ij\tilde{t}_{k}}^{n-} \mathcal{T}_{AA}^{nv-} + \overline{\mathcal{B}}_{ij\tilde{t}_{k}}^{n-} \mathcal{T}_{AB}^{nv-} \right) + \mathcal{B}_{jj\tilde{\mu}m}^{n-} \left(\overline{\mathcal{A}}_{ij\tilde{t}_{k}}^{n-} \mathcal{T}_{BA}^{nv-} + \overline{\mathcal{B}}_{ij\tilde{t}_{k}}^{n-} \mathcal{T}_{BB}^{n-} \right) \end{split}$$

where

$$\begin{split} \mathcal{T}_{AA}^{n\nu\pm} &= \left(\frac{1}{16\pi^2}\right) \frac{N_C Q_f}{4} \frac{m_{\mu}^2}{m_{\mu m}^2} \frac{m_i}{m_i \mp m_j} \\ &\times \left[\frac{1 - N_{f\tilde{t}_k} + l_{f\tilde{t}_k} N_{f\tilde{t}_k}^2}{(1 - N_{f\tilde{t}_k})^2 (N_{f\tilde{t}_k} - N_i)} - \frac{1 - N_i + l_i N_i^2}{(1 - N_i)^2 (N_{f\tilde{t}_k} - N_i)}\right] + (i \leftrightarrow j) \end{split}$$

and

$$N_{i/j} \equiv \frac{m_{i/j}^2}{m_{\tilde{\mu}m}^2}, \quad N_{f\tilde{f}_k} \equiv \frac{m_{f\tilde{f}_k}^2(w)}{m_{\tilde{\mu}m}^2}, \quad m_{f\tilde{f}_k}^2(w) \equiv \frac{m_{f}^2}{w} + \frac{m_{\tilde{f}_k}^2}{1-w}, \quad I_z \equiv \log N_z$$

MSSM fermion/sfermion two-loop corrections: sample analytic result

[arXiv:1311.1775][Fargnoli, Gnendiger, Paßehr, Stöckinger, S-K '13]

$$\begin{split} \Gamma_{ij\bar{t}_{k}}^{-,\mu}(\ell) &= \frac{1}{16\pi^{2}} e Q_{f'} \int_{0}^{1} \frac{\mathrm{d}w}{2} \Biggl\{ \left(\mathcal{A}_{ij\bar{t}_{k}}^{c\,+} - \mathcal{A}_{ij\bar{t}_{k}}^{c\,-} \gamma_{5} \right) \frac{\ell \not{g} \gamma^{\mu} - \ell q^{\mu} + \not{g} \ell^{\mu} - (\ell \cdot q) \gamma^{\mu}}{\mathcal{D}_{f'\bar{t}_{k}}(\ell)} \\ &+ \left(\mathcal{B}_{ij\bar{t}_{k}}^{c\,+} - \mathcal{B}_{ij\bar{t}_{k}}^{c\,-} \gamma_{5} \right) \frac{m_{f'}}{w} \frac{\not{g} \gamma^{\mu} - q^{\mu}}{\mathcal{D}_{f'\bar{t}_{k}}(\ell)} \Biggr\} \\ &- \frac{1}{16\pi^{2}} e Q_{\vec{\chi}} - \int_{0}^{1} \frac{\mathrm{d}w}{2} \Biggl\{ \left(\mathcal{A}_{ij\bar{t}_{k}}^{c\,+} - \mathcal{A}_{ij\bar{t}_{k}}^{c\,-} \gamma_{5} \right) w \Biggl[\left(\frac{1}{\epsilon} - \mathcal{L}(m_{\bar{t}_{k}}^{2}) \right) \gamma^{\mu} - 2 \frac{(\ell \cdot q) \ell}{\mathcal{D}_{f'\bar{t}_{k}}^{2}(\ell)} \\ &+ \frac{\not{g} \ell^{\mu} + \ell q^{\mu} + (\ell \cdot q) \gamma^{\mu} - 2 \ell \ell^{\mu} + \not{g}_{w}(\ell^{2}, m_{\bar{t}_{\ell}}^{2}, m_{\bar{t}_{k}}^{2}) \gamma^{\mu}/2}{\mathcal{D}_{f'\bar{t}_{k}}(\ell)} \Biggr] \\ &+ \left(\mathcal{B}_{ij\bar{t}_{k}}^{c\,+} - \mathcal{B}_{ij\bar{t}_{k}}^{c\,-} \gamma_{5} \right) m_{f'} \Biggl[- 2 \frac{(\ell \cdot q) \ell^{\mu}}{\mathcal{D}_{f'\bar{t}_{k}}^{2}(\ell)} + \frac{q^{\mu} - 2\ell^{\mu}}{\mathcal{D}_{f'\bar{t}_{k}}^{2}(\ell)} \Biggr] \Biggr\} \end{split}$$

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MSSM fermion/sfermion two-loop corrections: sample analytic result

[arXiv:1311.1775][Fargnoli, Gnendiger, Paßehr, Stöckinger, S-K '13]

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$$\begin{split} s^{cv}_{\mu,ij\tilde{t}_{k}} &= \int_{0}^{1} \mathrm{d} \mathsf{w} \bigg[\mathcal{A}^{c+}_{jj\tilde{\nu}\mu} \left(\mathcal{A}^{c+}_{ij\tilde{\nu}\mu} \mathcal{T}^{cv+}_{AA} + \mathcal{B}^{c+}_{ij\tilde{\nu}\mu} \mathcal{T}^{cv+}_{AB} \right) + \mathcal{B}^{c+}_{ji\tilde{\nu}\mu} \left(\mathcal{A}^{c+}_{ij\tilde{\nu}\mu} \mathcal{T}^{cv+}_{BA} + \mathcal{B}^{c+}_{ij\tilde{\nu}\mu} \mathcal{T}^{cv+}_{BB} \right) \\ &+ \mathcal{A}^{c-}_{ji\tilde{\nu}\mu} \left(\mathcal{A}^{c-}_{ij\tilde{\nu}\mu} \mathcal{T}^{cv-}_{AA} + \mathcal{B}^{c-}_{ij\tilde{\nu}\mu} \mathcal{T}^{cv-}_{AB} \right) + \mathcal{B}^{c-}_{ji\tilde{\nu}\mu} \left(\mathcal{A}^{c-}_{ij\tilde{\nu}\mu} \mathcal{T}^{cv-}_{BA} + \mathcal{B}^{c-}_{ij\tilde{\nu}\mu} \mathcal{T}^{cv-}_{BB} \right) \end{split}$$

$$\begin{aligned} \mathcal{T}_{AA}^{\text{cv2}}\left(C_{i}\right) &= \left[-12L(m_{\tilde{\nu}_{\mu}})-6l_{\tilde{l}_{k}}-3l_{i}+5-\frac{6}{C_{i}}\right]\frac{C_{i}\left(2-2C_{i}+3l_{i}C_{i}-l_{i}C_{i}^{2}\right)}{144\left(1-C_{i}\right)^{3}} \\ &+\frac{1+2l_{i}C_{i}-C_{i}^{2}}{12\left(1-C_{i}\right)^{3}} \end{aligned}$$

Non-decoupling behavior

[Physics Letters B 726 (2013)][arXiv:1309.0980][Fargnoli, Gnendiger, Paßehr, Stöckinger, S-K '13] BM1



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Non-decoupling behavior

[Physics Letters B 726 (2013)][arXiv:1309.0980][Fargnoli, Gnendiger, Paßehr, Stöckinger, S-K '13] BM4 can be largest 2L 0.15 contribution O(10%...30%)0.10 • $\mu = -160$, tan $\beta = 40$ 0.05 • $M_1 = 140 \text{GeV}$ 5 $m_{\tilde{\mu}_P} = 200 \text{GeV}$ 0.00 $M_2 = m_{\tilde{\mu}_1} = 2000 \text{GeV}$ -0.05 M_{U3} , $_{D3}$, $_{O3}$, $_{E3}$, $_{L3}$ M_{U}, D, Q $M_{O3}; M_{U3} = 1 TeV$ -0.10 $(tan \beta)$ 10^{3} 10^{4} 10^{5} 10^{6} photonic 2L(a) M [GeV]

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Leading logarithmic approximation

[http://iktp.tu-dresden.de/index.php?id=theory-software]

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$$\begin{split} \mathbf{a}_{\mu}^{\mathrm{LSUSY},\mathrm{M.I.}} &= \mathbf{a}_{\mu}^{\mathrm{LL}}(\tilde{W}-\tilde{H},\tilde{\nu}_{\mu}) + \mathbf{a}_{\mu}^{\mathrm{LL}}(\tilde{W}-\tilde{H},\tilde{\mu}_{L}) + \mathbf{a}_{\mu}^{\mathrm{LL}}(\tilde{B}-\tilde{H},\tilde{\mu}_{L}) \\ &\quad + \mathbf{a}_{\mu}^{\mathrm{LL}}(\tilde{B}-\tilde{H},\tilde{\mu}_{R}) + \mathbf{a}_{\mu}^{\mathrm{LL}}(\tilde{B},\tilde{\mu}_{L}-\tilde{\mu}_{R}) \\ \mathbf{a}_{\mu}^{\mathrm{LL}}(\tilde{W}/\tilde{B},\ldots) \propto g_{2/1}^{2}\mu \tan\beta \\ \mathbf{a}_{\mu}^{\mathrm{2L},\tilde{t}\tilde{f}\mathrm{LL}} &= \mathbf{a}_{\mu}^{\mathrm{LL}}(\tilde{W}-\tilde{H},\tilde{\nu}_{\mu}) \left(\Delta_{\tilde{g}2} + \Delta_{\tilde{H}} + \Delta_{\tilde{W}\tilde{H}} + \Delta_{t_{\beta}} + 0.015\right), \\ &\quad + \mathbf{a}_{\mu}^{\mathrm{LL}}(\tilde{W}-\tilde{H},\tilde{\mu}_{L}) \left(\Delta_{g_{2}} + \Delta_{\tilde{H}} + \Delta_{\tilde{W}\tilde{H}} + \Delta_{t_{\beta}} + 0.015\right), \\ &\quad + \mathbf{a}_{\mu}^{\mathrm{LL}}(\tilde{B}-\tilde{H},\tilde{\mu}_{L}) \left(\Delta_{g_{1}} + \Delta_{\tilde{H}} + \Delta_{\tilde{B}\tilde{H}} + \Delta_{t_{\beta}} + 0.015\right), \\ &\quad + \mathbf{a}_{\mu}^{\mathrm{LL}}(\tilde{B}-\tilde{H},\tilde{\mu}_{R}) \left(\Delta_{g_{1}} + \Delta_{\tilde{H}} + \Delta_{\tilde{B}\tilde{H}} + \Delta_{t_{\beta}} + 0.04\right), \\ &\quad + \mathbf{a}_{\mu}^{\mathrm{LL}}(\tilde{B},\tilde{\mu}_{L}-\tilde{\mu}_{R}) \left(\Delta_{g_{1}} + \Delta_{t_{\beta}} + 0.03\right) \\ \Delta_{\tilde{H}} &= \frac{1}{16\pi^{2}} \frac{1}{2} \left(3y_{t}^{2}\log\frac{M_{U3}}{m_{\mathrm{SUSY}}} + 3y_{b}^{2}\log\frac{M_{D3}}{m_{\mathrm{SUSY}}} \\ &\quad + 3(y_{t}^{2} + y_{b}^{2})\log\frac{M_{Q3}}{m_{\mathrm{SUSY}}} + y_{\tau}^{2}\log\frac{M_{L3}}{m_{\mathrm{SUSY}}} + y_{\tau}^{2}\log\frac{M_{L3}}{m_{\mathrm{SUSY}}}\right) \end{split}$$

Leading logarithmic approximation



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Summary

- a_{μ}^{SM} : a_{μ}^{EW} contributions depending on M_H \rightarrow this uncertainty has been eliminated
- a_{μ} still viable, complementary constraint on SUSY
- $a_{\mu}^{2L,\gamma}$: -(7...9)% corrections to $a_{\mu}^{1L,SUSY}$
- fermion/sfermion two-loop contributions:
 - first full calculation of a_{μ}^{SUSY} two-loop contributions: 5-heavy-mass-scale diagrams
 - leading logarithmic approximation: invalid when $\tan\beta$ or M_{SUSY} is very small
 - ln(m_f) enhanced, non-decoupling behavior for (very) heavy squarks: 10...30% corrections to a^{1L,SUSY}_µ possible