# Wino Dark Matter Revisited

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Why is the wino dark matter attractive?
Phenomenology of the wino dark matter
Summary & Discussions

Why is the wino DM attractive?

### Wino DM = SU(2)<sub>L</sub>-triplet Majorana fermion





- Residual Z<sub>2</sub> symmetry of  $U(1)_{B-L}$  -Neutrino masses & mixings + BAU  $\downarrow$ Three heavy R-handed neutrinos.  $\downarrow$ Gauged U(1)<sub>B-L</sub> ( $\because$  anomaly cancel)

#### Why U(1)<sub>B-L</sub> works?

- $U(1)_{B-L}$  can be spontaneously broken by a VEV having B-L charge 2.
- The SM involves only B-L odd fermions and B-L even (zero) bosons,

A New fermion (boson) with a even (odd) B-L charge becomes stable! Particle (either fermion or boson) with a fractional B-L charge is OK.

Wino DM is nothing but a fermion which is singlet under  $U(1)_{B-L}$  !

Wino DM is attractive from the viewpoint of DM stability!

Why is the wino DM attractive?

#### Wino DM = SU(2)<sub>L</sub>-triplet Majorana fermion

Recent LHC results may indicate that BSM scale is much higher than 1TeV, and the Lagrangian below O(1)TeV is minimally given by SM + DM system.

SU(2) →	singlet	doublet	triplet
Scalar	1	3	5
Fermion	2	4	6

2/8

is severely limited by direct detection experiments. ( : (φ<sub>DM</sub>)<sup>2</sup> |H|<sup>2</sup>)
does not have any renormalizable interactions with SM particles.
must have a hypercharge to have a neutral component, resulting in a complex scalar DM and severely constrained by direct detections. (If we introduce higher dimensional operators to split the complex scalar into two real ones, the direct detection limit can be evaded.)
has the same situation as the case of 3.

(5) can be a real scalar (representation), but severely constrained by direct detection experiments because of (φ<sub>DM</sub>)<sup>2</sup> |H|<sup>2</sup> type interaction.
(6) can be a Majorana fermion. No worry about direct detection limits.

Wino DM is attractive from the viewpoint of minimality!

Why is the wino DM attractive?

### Wino DM = SU(2)<sub>L</sub>-triplet Majorana fermion

In the framework of the MSSM, the discovery of the  $\sim 126$ GeV higgs boson may indicate that the typical scale of sfermion masses is 0(10-1000)TeV. High-scale SUSY scenarioPGM[Ibe, Moroi, Yanagida, PLB644]mSplit[A. Arvanitakia, et.al., JHEP1302]Spread[L. Hall, et.al., JHEP1031(2013)]KLM[E. Dudas, et.al., EPJ C73 (2013)]



Wino DM is attractive from the viewpoint of the MSSM!

0.3

0.2

0

0.1 WMAP

4/8

Thermal relic abundance,  $\Omega_{\rm DM}h^2$ 

2

m (TeV)

Perturbative

3

Non-perturbative

[J.Hisano, S.M. M.Nagai, O.Saito,

**M.Senami** (2007)]

Upper bound of m<sub>wino</sub> from cosmology

#### Two contributions to $\Omega_{DM}h^2$

**1**. Thermal contribution  $\Omega_{TH} h^2 \rightarrow$ **2**. Non-thermal contribution  $\Omega_{NT} h^2$ 

(e.g., via gravitino decays, etc.) [J.Feng, T.Moroi, L.Randall, M.Strassler, S.Su (1999)]

$$(\Omega_{TH} + \Omega_{NT}) h^2 = 0.11$$
  
 $\rightarrow m_{wino} < 2.8 TeV$ 

#### Lower bound of m<sub>wino</sub> from LHC



5/8

 $\sim$  Dark Matter detections to put robust limits  $\sim$ 

Detections	Meth	od	Experiments	Comments
Direct			XENON100/LUX	$\sigma_{\chi N} \sim 1.5 \times 10^{-47}  \mathrm{cm^2}$ , so that it is difficult to detect it in near future.
Indirect CR		19 24	PAMELA/AMS-02	Propagation uncertainty is large.
Indirect $v$			IceCube/DeepCore	Detection efficiency is small.
<b>Indirect</b> γ	dSph	line	Fermi-LAT/HESS	Signal flux is small to be detected.
	CGR	line		Uncertainty of DM profile is large, but large DM density in this region,
	Cluck	cont.		" + Difficulty of BG estimation.
	Clust.	cont.	Surger 1	Good, but dSph observation is better.
2. 动脉系统	Diff.	line		Signal flux is small to be detected.
		cont.		Good, but dSph observation is better.

Indirect detection of the wino DM has an advantage, for its annihilation cross section is boosted by the Sommerfeld effect. [J. Hisano, S. M. M. Nojiri, PRL92 (2004)]

### Flux from galaxy satellites $\frac{d\mathcal{F}}{dE} = \frac{\langle \sigma v \rangle}{8\pi m^2} \sum_{f} \operatorname{Br}(f) \frac{dN_f}{dE} \int_{\delta\Omega} d\Omega \int_{l.o.s.} dl \,\rho^2(\vec{r})$

CL. dSphs: Ursa Minor, Draco, … (DM profile is well estimated.) UF. dSphs: Ursa Majorll, Seque1, … (Kinematical data is still poor.)

Robust limit (Right-upper fig.)
→ 0.32TeV < m<sub>wino</sub> < 2.25TeV</li>
2.44TeV < m<sub>wino</sub>

• Future limit (Right-lower fig.) It depends how small the error of J factor can be in each dSph. If  $\Delta(\log_{10}J) = 0.1$ , all region covered.

#### Current limit



6/8

Mwino

0.1TeV

J-factor!

**1TeV** 



# **Summary & Discussions**

- According to current results at LHC on the higgs discovery and new physics searches, the neutral wino seems to be attractive! The wino dark matter is indeed favored from viewpoints of its stability, minimality, and MSSM (in the high-scale SUSY scenario).
- Robust upper and lower limits on the wino mass are obtained by cosmology and LHC, 0.27TeV <  $m_{wino}$  < 2.8TeV. DM signals at the direct detection experiments are very weak, while DM signals at indirect detection experiments are good. From continuum  $\gamma$ -ray observations from dSphs and line  $\gamma$ -ray from CGS put limits as 0.32TeV < m<sub>wino</sub> < 2.1TeV & 2.6TeV < m<sub>wino</sub> < 2.8TeV.
- In near future,  $\gamma$ -ray observations still play important roles to test the wino dark matter, because its systematic error can be well controlled. Continuum  $\gamma$ -ray from dSphs can be completely test the dark matter if errors on J-factor are reduced. About line  $\gamma$ -ray from CGR, observing  $\gamma$ -ray from regions around the exact center will be also important for cored DM profiles.